

Anomalous Diffusion and Basic Theorems on Statistical Mechanics

Fernando Albuquerque de Oliveira

Instituto de Física
and
International Center for Condensed Matter Physics
Universidade de Brasília



- 1 Motivation
- 2 Brownian Motion
- 3 Memory
- 4 Irreversibility and ballistic Motion
- 5 Ergodicity and Khinchin Theorem
- 6 Conclusions

Basic Theorems

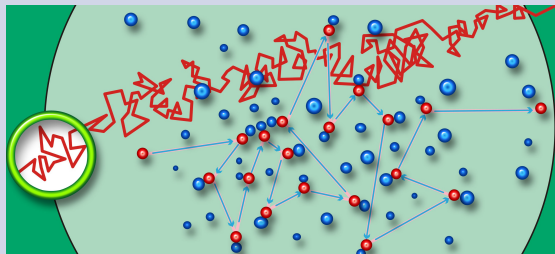


- How a system approaches equilibrium?
- Ergodic Hypothesis.
- Khinchin Theorem.
- Fluctuation Dissipation Theorem.
- The Fluctuation Dissipation Theorem.
- We can discuss all those concepts in diffusion.

General Concepts



Brownian Motion



Einstein Relation

$$\lim_{t \rightarrow \infty} \langle x^2(t) \rangle \sim t^\alpha$$

Diffusion Exponent

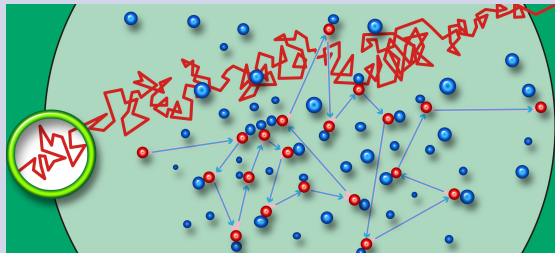
- $\alpha < 1$, Subdiffusion;
- $\alpha = 1$, Normal Diffusion;
- $\alpha > 1$, Superdiffusion.

$$\lim_{t \rightarrow \infty} \langle x^2(t) \rangle \sim 2Dt, \text{ with } D = \frac{k_B T}{m\gamma}$$

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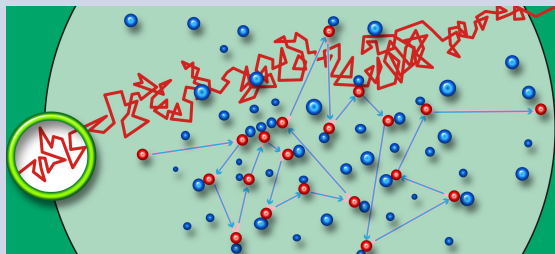
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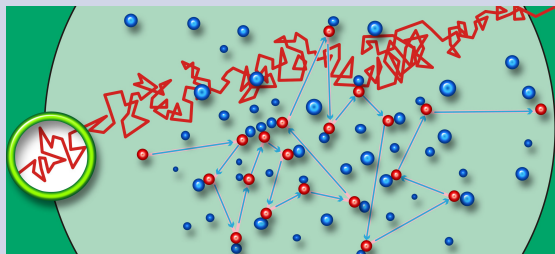
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Langevin's Equation



Langevin (1908)

$$\frac{d}{dt}p(t) = -\gamma p(t) + \xi(t)$$

Properties

- The noise $\xi(t)$ fulfils $\langle \xi(t) \rangle = 0$ e $\langle p(0)\xi(t) \rangle = 0$.
- Fluctuation-Dissipation Theorem:

$$\langle \xi(t)\xi(t') \rangle = 2m\gamma k_B T \delta(t-t')$$

Diffusion Constant

$$D = \int_0^\infty C_v(t) dt = \int_0^\infty \langle v(t)v(0) \rangle dt = \frac{k_B T}{m\gamma}$$

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Generalized Langevin's Equation



Mori equation's (1964)

$$m \frac{dv(t)}{dt} = -m \int_0^t \Gamma(t-t') v(t') dt' + \xi(t)$$

Properties

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- The friction is time (frequency) dependent.
- Generalized Fluctuation-Dissipation Theorem:

$$\langle \xi(t) \xi(t') \rangle = \langle (mv)^2 \rangle_{eq} \Gamma(t-t')$$

Solution

$$v(t) = v(0)R(t) + \int_0^t \xi(t') R(t-t') dt'$$

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Anomalous diffusion



Correlation function $C_V(t)$

$$\frac{dC_V(t)}{dt} = - \int_0^t \Gamma(t-t') C_V(t') dt'$$

- If $\langle x^2(t) \rangle = Dt$ and $\langle x^2(t) \rangle \sim t^\alpha$ for $\alpha \neq 1$.
- D must be a function of t

$$\lim_{t \rightarrow \infty} D(t) = \lim_{t \rightarrow \infty} \int_0^t C_V(t') dt' = \lim_{z \rightarrow 0} \int_0^\infty C_V(t) \exp(-zt) = \lim_{z \rightarrow 0} \tilde{C}_V(z)$$

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Classification of diffusion



- Consider $\lim_{z \rightarrow 0} \tilde{\Gamma}(z) \sim z^\nu$
- $-1 < \nu < 1$

$$\lim_{t \rightarrow \infty} D(t) = \lim_{z \rightarrow 0} \tilde{C}_\nu(z) = \lim_{z \rightarrow 0} \frac{C_\nu(0)}{z + \tilde{\Gamma}(z)} \sim \lim_{z \rightarrow 0} \tilde{\Gamma}(z)^{-1} \sim t^\nu$$

$$\langle x^2(t) \rangle = 2D(t)t \sim t^{\nu+1} \sim t^\alpha$$

$$\alpha = \nu + 1$$



F. A. Oliveira et al , *Phy. Rev. Lett.* 86, 5839 (2001).



R. Morgado et al , *Phy. Rev. Lett.* 89, 100601 (2002).

From the noise

$$\xi(t) = \int_0^\infty \sqrt{\frac{\rho_r(\omega)}{2mk_B T}} \cos[\omega t + \phi(\omega)] d\omega$$

- $\rho_r(\omega)$ is the spectral density
- ϕ are random numbers $0 < \phi < 2\pi$

Fluctuation Dissipation

$$\langle \xi(t) \xi(t') \rangle = \langle (mv)^2 \rangle_{eq} \Gamma(t-t')$$

$$\Gamma(t) = \int_0^\infty \rho_r(\omega) \cos(\omega t) d\omega$$

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Memory and spectral density

$$\Gamma(t) = \int_0^\infty \rho_r(\omega) \cos(\omega t) d\omega$$

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- The spectral density determines the friction

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- The spectral density determines the friction

Irreversibility

$$\lim_{t \rightarrow \infty} R(t) = 0$$

$$D = \lim_{t \rightarrow \infty} \frac{\langle x^2(t) \rangle}{2t} \Rightarrow \lim_{z \rightarrow 0} \tilde{\Gamma}(z) = bz^\alpha \text{ with } \alpha = \nu + 1$$

$$\lim_{t \rightarrow \infty} R(t) = \lim_{z \rightarrow 0} R(z) = \lim_{z \rightarrow 0} \frac{z}{\tilde{\Gamma}(z)} = \lim_{z \rightarrow 0} \left(1 - \frac{\tilde{\Gamma}(z)}{z}\right) \quad \text{now } \frac{\tilde{\Gamma}(z)}{z} = bz^{\nu-1}$$

Anomalous Diffusion

- $-1 < \nu < 1$ ($0 < \alpha < 2$)

- $\lim_{z \rightarrow 0} \frac{\tilde{\Gamma}(z)}{z} \rightarrow \infty$

- $R(t \rightarrow \infty) = 0$

Ballistic Diffusion

- $\nu = 1$ ($\alpha = 2$)

- $R(t \rightarrow \infty) = \frac{1}{1+b} = \kappa \neq 0$

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Anomalous Diffusion

$$\blacksquare -1 < \nu < 1 \quad (0 < \alpha < 2)$$

$$\blacksquare \lim_{z \rightarrow 0} \frac{\tilde{\Gamma}(z)}{z} \rightarrow \infty$$

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$$\lim_{t \rightarrow \infty} R(t) = 0 \Leftrightarrow \text{Ergodicity holds}$$

Lee conjecture's (prl 2007)

$p(t)$ is ergodic if and only if,

$$0 < W = \int_0^\infty R(t) dt < \infty$$

Our calculations

$$\langle p^2(t) \rangle = \langle p^2 \rangle_{eq} + R^2(t) \left[\langle p^2(0) \rangle - \langle p^2 \rangle_{eq} \right]$$

- The Khinchin Theorem holds for all kind of diffusion



L. C. Lapas, R. Morgado, M. H. Vainstein, J. M. Rubí, and F. A. Oliveira, *Phy. Rev. Lett.* 101, 230602 (2008).

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Fluctuation-Dissipation Theorem

- Ergodicity does work for anomalous diffusion
- Ergodicity does not holds for ballistic diffusion,
- The same for the fluctuation-dissipation theorem.



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Friction and memory

$$\gamma = \lim_{z \rightarrow 0} \tilde{\Gamma}(z) \sim z^{\nu} \rightarrow \begin{cases} 0 & , \text{ Superdiffusion } (\nu > 0) \\ \text{constant} & , \text{ Normal dif. } (\nu = 0) \\ \infty & , \text{ Subdiffusion } (\nu < 0) \end{cases}$$

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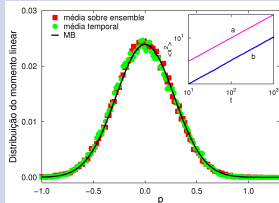
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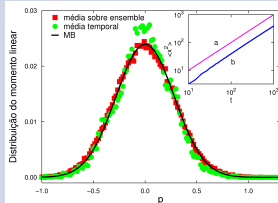
Ergodicity in Anomalous Diffusion



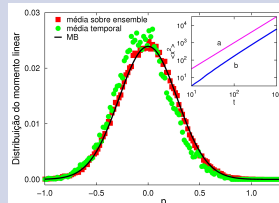
Subdiffusion $\alpha = 0.5$



Normal Diffusion



Superdiffusion $\alpha = 1.5$



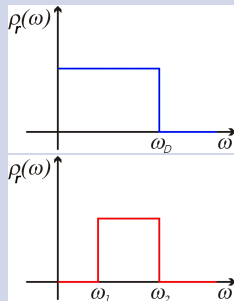
Noise spectral density:

$$\rho_r(\omega) = \begin{cases} \frac{2\gamma_0}{\pi} & , 0 < \omega < \omega_D \\ 0 & , \text{otherwise} \end{cases}$$

Process Ornstein-Zernike:

$$\rho_{DB}(\omega) = \rho_{\omega_2}(\omega) - \rho_{\omega_1}(\omega)$$

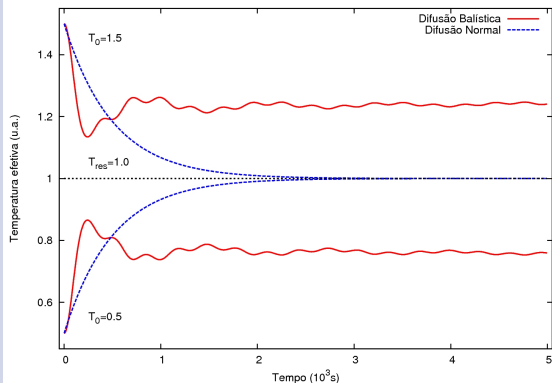
Normal and ballistic diffusion



Memory Function

$$\Gamma(t) = \frac{2\gamma_0}{\pi} \left[\frac{\sin(\omega_2 t)}{t} - \frac{\sin(\omega_1 t)}{t} \right]$$

Effective Temperature



Entropy



■ Variation of entropy:

$$\Delta S = \int_{T_0}^{T_{ef}} c_v(T') \left(\frac{1}{T'} - \frac{1}{T_{res}} \right) dT' \geq 0.$$

■ Residual entropy:

$$\Delta S = \int_{T_{res}}^{T_{ef}} c_v(T') \left(\frac{1}{T'} - \frac{1}{T_{res}} \right) dT' \geq 0$$

■ final entropy:

$$S = S_{max} - \Delta S'$$

■ Gibbs entropy:

Maximum of entropy $\iff$ Gaussian distribution

$$S(t) = -k_B \int \rho(p, t) \ln \frac{\rho(p, t)}{\rho_{eq}(p)} dp$$

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Non Gaussian Behavior



Third momentum

$$\langle p^3(t) \rangle = \langle p^3(0) \rangle R^3(t) + 3 \langle p(0) \rangle \langle p^2 \rangle_{eq} [1 - R^2(t)] R(t)$$

Forth momentum

$$\begin{aligned} \langle p^4(t) \rangle = & \langle p^4(0) \rangle R^4(t) + 3 \langle p^2 \rangle_{eq}^2 [1 - R^2(t)]^2 + \\ & 6 \langle p^2(0) \rangle \langle p^2 \rangle_{eq} [1 - R^2(t)] R^2(t) \end{aligned}$$

Skewness

$$\zeta(t) = \left[\frac{\sigma_p(0)}{\sigma_p(t)} \right]^3 \zeta(0) R^3(t)$$

Non Gaussian factor

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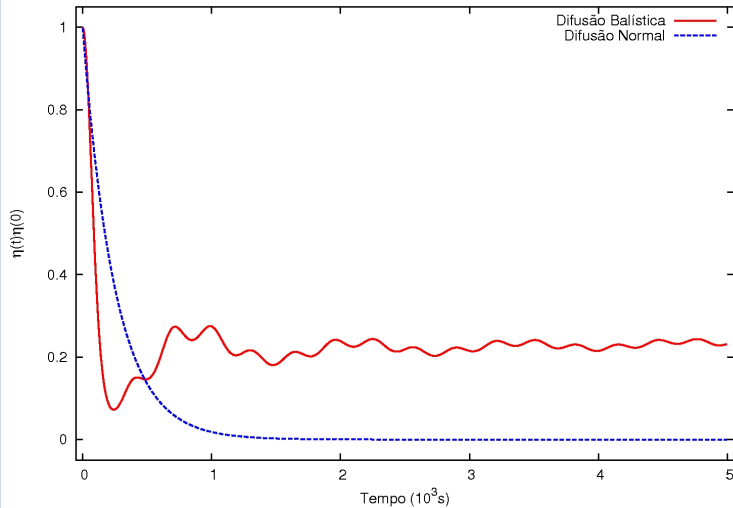
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Conclusions

- Anomalous diffusion $\lim_{t \rightarrow \infty} \langle x^2(t) \rangle \sim t^\alpha$ holds for $0 \leq \alpha \leq 2$.
- From the memory $\lim_{z \rightarrow 0} \tilde{\Gamma}(z) \sim z^\nu$ we get $\alpha = \nu + 1$.
- Second law is valid for all diffusive regimes.
- All process ends up in a Gaussian, except the ballistic one;
- The ballistic evolves towards a Gaussian without reach it.
- The temperature evolves towards the reservoir temperature.
- The symmetry of the distribution is preserved;
- Ergodicity works for all diffusive regimes except for ballistic.
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